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## LETTER TO THE EDITOR

# On the icosahedral equation and the locus of zeros for the grand partition function of the hard-hexagon model 

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#### Abstract

The Kleinian theory of the icosahedral equation is used to investigate a recent conjecture on the locus of zeros for the grand partition function of the hard-hexagon lattice gas model.


Baxter $(1980,1981)$ has shown that in the thermodynamic limit the grand partition function per site $\Xi$ of the hard-hexagon lattice gas model has a parametric representation in the ordered regime given by

$$
\begin{align*}
& \Xi=x^{-1 / 3} \frac{G^{3}(x)}{H^{2}(x)} \prod_{n=1}^{\infty} \frac{\left(1-x^{3 n-2}\right)\left(1-x^{3 n-1}\right)\left(1-x^{5 n}\right)^{2}}{\left(1-x^{3 n}\right)^{2}}  \tag{1}\\
& z^{-1}=x[H(x) / G(x)]^{5} \tag{2}
\end{align*}
$$

where $z$ is the activity of the gas, and

$$
\begin{align*}
& G(x)=\prod_{n=1}^{\infty}\left[\left(1-x^{5 n-4}\right)\left(1-x^{5 n-1}\right)\right]^{-1}  \tag{3}\\
& H(x)=\prod_{n=1}^{\infty}\left[\left(1-x^{5 n-3}\right)\left(1-x^{5 n-2}\right)\right]^{-1} \tag{4}
\end{align*}
$$

and $0<x<1$. When $x \rightarrow 1$ - the model displays an order-disorder transition with a critical activity value

$$
\begin{equation*}
z_{\mathrm{c}}=\left[\frac{1}{2}(1+\sqrt{ } 5)\right]^{5} . \tag{5}
\end{equation*}
$$

Similar results were also obtained by Baxter for the disordered regime $0 \leqslant z<z_{\mathrm{c}}$.
Recently, Joyce (1988) has used the theory of modular functions to eliminate the non-physical parameter $x$ from equations (1) and (2). It was found that the grand partition function is an algebraic function $\Xi\left(z^{\prime}\right)$ of the reciprocal activity $z^{\prime}=z^{-1}$ which satisfies the polynomial equation

$$
\begin{align*}
\left(z^{\prime}\right)^{2} \Omega_{1}^{10}\left(z^{\prime}\right) y^{4}- & \Omega_{3}\left(z^{\prime}\right)\left[1458 z^{\prime} \Omega_{1}^{5}\left(z^{\prime}\right)+\Omega_{3}^{2}\left(z^{\prime}\right)\right] y^{3} \\
& -3^{10}\left[2430 z^{\prime} \Omega_{1}^{5}\left(z^{\prime}\right)+\Omega_{3}^{2}\left(z^{\prime}\right)\right] y^{2}-3^{19} \Omega_{3}\left(z^{\prime}\right) y-3^{27}=0 \tag{6}
\end{align*}
$$

where

$$
\begin{align*}
& y=\Xi^{6}  \tag{7}\\
& \Omega_{1}\left(z^{\prime}\right)=1-11 z^{\prime}-\left(z^{\prime}\right)^{2}  \tag{8}\\
& \Omega_{3}\left(z^{\prime}\right)=1-522 z^{\prime}-10005\left(z^{\prime}\right)^{2}-10005\left(z^{\prime}\right)^{4}+522\left(z^{\prime}\right)^{5}+\left(z^{\prime}\right)^{6} . \tag{9}
\end{align*}
$$

Wood et al (1989) have investigated the resolvent properties of the polynomial equation (6) and have conjectured that part of the limiting locus of grand partition function zeros in the $z$ plane for the hard-hexagon model can be generated by considering the solutions $z=z(w)$ of the rational algebraic equation

$$
\begin{equation*}
w=z^{\prime} \Omega_{1}^{5}\left(z^{\prime}\right) / \Omega_{3}^{2}\left(z^{\prime}\right) \tag{10}
\end{equation*}
$$

where $w$ is a real parameter.
The main purpose in this letter is to investigate the mathematical properties of the solutions $z=z(w)$ of equation (10). In particular, it will be shown that the algebraic curves generated by the solutions $z=z(w)$ have simple rational parametric representations. We begin the analysis by considering the following basic modular functions (Klein and Fricke 1890, p 154, 1892, p 383):

$$
\begin{align*}
& J(\tau)=(1728 x)^{-1}\left[1+240 \sum_{n=1}^{\infty} n^{3} x^{n}\left(1-x^{n}\right)^{-1}\right]^{3} \prod_{n=1}^{\infty}\left(1-x^{n}\right)^{-24}  \tag{11}\\
& \zeta(\tau)=x^{1 / 5} H(x) / G(x) \tag{12}
\end{align*}
$$

where

$$
\begin{equation*}
x=\exp (2 \pi \mathrm{i} \tau) \tag{13}
\end{equation*}
$$

and the variable $\tau$ lies in the upper half-plane $\operatorname{Im}(\tau)>0$. The function $J(\tau)$ is the fundamental hauptmodul for the full modular group $\Gamma$, and $\zeta(\tau)$ is the hauptmodul for the principal congruence subgroup $\Gamma(5)$. Because the function $J(\tau)$ is a modular invariant for all the transformations belonging to $\Gamma(5)$ it is possible to write $J(\tau)$ as a rational function of $\zeta(\tau)$. In fact Klein and Fricke (1890, p 105) have proved that this rational relation can be expressed in the two equivalent forms:

$$
\begin{align*}
& 1728 J=\Omega_{2}^{3}\left(\zeta^{5}\right) / \zeta^{5} \Omega_{1}^{5}\left(\zeta^{5}\right)  \tag{14}\\
& 1728(J-1)=\Omega_{3}^{2}\left(\zeta^{5}\right) / \zeta^{5} \Omega_{1}^{5}\left(\zeta^{5}\right) \tag{15}
\end{align*}
$$

where

$$
\begin{equation*}
\Omega_{2}\left(\zeta^{5}\right)=1+228 \zeta^{5}+494 \zeta^{10}-228 \zeta^{15}+\zeta^{20} \tag{16}
\end{equation*}
$$

and the polynomials $\Omega_{1}$ and $\Omega_{3}$ are defined in equations (8) and (9) respectively.
Equation (14) defines an algebraic inverse function $\zeta(J)$ which consists of 60 function elements, and has a Galois group which is isomorphic with the icosahedral rotation group (Klein 1913). It is also known that the zeros of the polynomials $\Omega_{j}\left(\zeta^{5}\right)$ $(j=1,2,3)$ are closely connected with the geometrical properties of the icosahedron. Because of these remarkable properties equations (14) and (15) are, not surprisingly, referred to as the icosahedral equations. If we now make the identifications

$$
\begin{align*}
& z^{\prime}=z^{-1} \equiv \zeta^{5}  \tag{17}\\
& w^{-1} \equiv 1728(J-1) \tag{18}
\end{align*}
$$

we see that the basic equation (10) becomes the icosahedral equation (15) with $J$ real.
Unfortunately, the icosahedral equation (15) is not solvable in terms of radicals because the Galois group for the equation is simple. However, it is possible to express the inverse function $\zeta(J)$ in terms of hypergeometric series by using the work of Schwarz (1873). In this manner we find that one solution of the equation is

$$
\begin{equation*}
\zeta_{0}(J)=(12)^{-3 / 5} \operatorname{sgn}(J)|J|^{-1 / 5} F_{1} / F_{2} \tag{19}
\end{equation*}
$$

where

$$
\begin{align*}
& F_{1} \equiv{ }_{2} F_{1}\left(\frac{11}{60}, \frac{31}{60} ; \frac{6}{5} ; 1 / J\right)  \tag{20}\\
& F_{2} \equiv{ }_{2} F_{1}\left(-\frac{1}{60}, \frac{19}{60} ; \frac{4}{5} ; 1 / J\right) \tag{21}
\end{align*}
$$

and $1 \leqslant|J|<\infty$, with $J$ real. The application of the transformation formula (Erdélyi et al 1953, p 105)

$$
\begin{equation*}
{ }_{2} F_{1}(a, b ; c ; z)=(1-z)^{-a}{ }_{2} F_{1}[a, c-b ; c ; z /(z-1)] \tag{22}
\end{equation*}
$$

to equation (19) yields the alternative expression

$$
\begin{equation*}
\zeta_{0}(J)=(12)^{-3 / 5} \operatorname{sgn}(J-1)|J-1|^{-1 / 5} F_{3} / F_{4} \tag{23}
\end{equation*}
$$

where

$$
\begin{align*}
& F_{3} \equiv{ }_{2} F_{1}\left(\frac{11}{60}, \frac{41}{60} ; \frac{6}{5} ; 1 /(1-J)\right)  \tag{24}\\
& F_{4} \equiv{ }_{2} F_{1}\left(-\frac{1}{60}, \frac{29}{60} ; \frac{4}{5} ; 1 /(1-J)\right) \tag{25}
\end{align*}
$$

and $-\infty<J \leqslant 0$ or $J>2$. It is clear that equation (23) is particularly useful for the case $J \leqslant 0$.

We can determine the behaviour of $\zeta_{0}(J)$ in the neighbourhood of the branch point $J=1$ by applying standard analytic continuation formulae (Erdélyi et al 1953, pp 105, 108) to equation (19). The final result is

$$
\begin{equation*}
\zeta_{0}(J)=A_{1}\left[F_{5}-A_{2}(J-1)^{1 / 2} F_{6}\right] /\left[F_{5}+A_{3}(J-1)^{1 / 2} F_{6}\right] \tag{26}
\end{equation*}
$$

where

$$
\begin{align*}
& F_{5} \equiv{ }_{2} F_{1}\left(-\frac{1}{60}, \frac{11}{60} ; \frac{1}{2} ; 1-J\right)  \tag{27}\\
& F_{6} \equiv{ }_{2} F_{1}\left(\frac{29}{60}, \frac{41}{60} ; \frac{3}{2} ; 1-J\right)  \tag{28}\\
& A_{1}=2^{-1}\left[-(1+\sqrt{ } 5)+\sqrt{ } 2(5+\sqrt{ } 5)^{1 / 2}\right]  \tag{29}\\
& A_{2}=2^{-2} 3^{-1 / 2} 5^{-3 / 4}\left[(1+\sqrt{ } 5)+\sqrt{ } 2(5+\sqrt{ } 5)^{1 / 2}\right]  \tag{30}\\
& A_{3}=2^{-2} 3^{-1 / 2} 5^{-3 / 4}\left[-(1+\sqrt{ } 5)+\sqrt{ } 2(5+\sqrt{ } 5)^{1 / 2}\right] \tag{31}
\end{align*}
$$

and $1 \leqslant J \leqslant 2$. A similar procedure can be used to establish the behaviour of $\zeta_{0}(J)$ in the neighbourhood of the branch point $J=0$. It is found that

$$
\begin{equation*}
\zeta_{0}(J)=-A_{4}\left[F_{7}+A_{5} \operatorname{sgn}(J)|J|^{1 / 3} F_{8}\right] /\left[F_{7}-A_{6} \operatorname{sgn}(J)|J|^{1 / 3} F_{8}\right] \tag{32}
\end{equation*}
$$

where

$$
\begin{align*}
& F_{7} \equiv{ }_{2} F_{1}\left(-\frac{1}{60}, \frac{11}{60} ; \frac{2}{3} ; J\right)  \tag{33}\\
& F_{8} \equiv{ }_{2} F_{1}\left(\frac{19}{60}, \frac{31}{60} ; \frac{4}{3} ; J\right)  \tag{34}\\
& A_{4}=2^{-2}\left[-(3+\sqrt{ } 5)+\sqrt{ } 6(5+\sqrt{ } 5)^{1 / 2}\right]  \tag{35}\\
& A_{5}=2^{-3} 5^{-5 / 6}\left[(3+\sqrt{ } 5)+\sqrt{ } 6(5+\sqrt{ } 5)^{1 / 2}\right]  \tag{36}\\
& A_{6}=2^{-3} 5^{-5 / 6}\left[-(3+\sqrt{ } 5)+\sqrt{ } 6(5+\sqrt{ } 5)^{1 / 2}\right] \tag{37}
\end{align*}
$$

and $-1 \leqslant J<1$. Finally, we note that the behaviour of the solution (32) in the neighbourhood of $J=1$ is given by the formula

$$
\begin{equation*}
\zeta_{0}(J)=-A_{7}\left[F_{5}-A_{8}(1-J)^{1 / 2} F_{6}\right] /\left[F_{5}+A_{9}(1-J)^{1 / 2} F_{6}\right] \tag{38}
\end{equation*}
$$

where

$$
\begin{align*}
& A_{7}=2^{-1}\left[(1-\sqrt{ } 5)+\sqrt{ } 2(5-\sqrt{ } 5)^{1 / 2}\right]  \tag{39}\\
& A_{8}=2^{-2} 3^{-1 / 2} 5^{-3 / 4}\left[-(1-\sqrt{ } 5)+\sqrt{ } 2(5-\sqrt{ } 5)^{1 / 2}\right]  \tag{40}\\
& A_{9}=2^{-2} 3^{-1 / 2} 5^{-3 / 4}\left[(1-\sqrt{ } 5)+\sqrt{ } 2(5-\sqrt{ } 5)^{1 / 2}\right] \tag{41}
\end{align*}
$$

and $0 \leqslant J<1$. We have now constructed one real root $\zeta_{0}(J)$ of the icosahedral equation for each real value of $J$. It should be pointed out, however, that the various formulae given for $\zeta_{0}(J)$ do not all represent the same branch of the algebraic function $\zeta(J)$.

Klein (1922) has shown that if we have found one root $\zeta_{0}=\zeta_{0}(J)$ of the icosahedral equation then all the other roots of the equation can be expressed as fractional linear transformations of $\zeta_{0}$. The detailed results are

$$
\begin{align*}
\zeta=\zeta(J) & =\varepsilon^{\mu} \zeta_{0}  \tag{42}\\
& =-\varepsilon^{\mu} / \zeta_{0}  \tag{43}\\
& =\varepsilon^{\mu}\left[\frac{\left(\varepsilon+\varepsilon^{4}\right) \zeta_{0}+\varepsilon^{\nu}}{\zeta_{0}-\varepsilon^{\nu}\left(\varepsilon+\varepsilon^{4}\right)}\right]  \tag{44}\\
& =-\varepsilon^{\mu}\left[\frac{\zeta_{0}-\varepsilon^{\nu}\left(\varepsilon+\varepsilon^{4}\right)}{\left(\varepsilon+\varepsilon^{4}\right) \zeta_{0}+\varepsilon^{\nu}}\right] \tag{45}
\end{align*}
$$

where

$$
\begin{equation*}
\varepsilon=\exp (2 \pi \mathrm{i} / 5) \tag{46}
\end{equation*}
$$

$\mu, \nu=0,1,2,3,4$ and $\zeta_{0}=\zeta_{0}(J)$. From these formulae we can now readily derive the following explicit expressions for the 12 roots $z=z(J)$ of the basic equation (10) with $J$ real:

$$
\begin{equation*}
z_{0, \pm m}=z_{\mathrm{c}}\left[\frac{\frac{1}{2}(\sqrt{ } 5-1) \zeta_{0}+\varepsilon^{ \pm m}}{-\frac{1}{2}(\sqrt{ } 5+1) \zeta_{0}+\varepsilon^{ \pm m}}\right]^{5} \tag{47}
\end{equation*}
$$

$z_{1, \pm m}=-1 / z_{0, \pm m}, z_{2}=\zeta_{0}^{-5}, z_{3}=-1 / z_{2}, z_{4}=z_{0,0}$ and $z_{5}=-1 / z_{4}$, where $m=1,2$ and $\zeta_{0}=\zeta_{0}(J)$ is the real root of the icosahedral equation. If the complex roots $z_{0, \pm 1}(J)$ and $z_{1, \pm 1}(J)$ are plotted for real values of $J$ in the $z$ plane one obtains a locus which is a complicated closed curve $C_{1}$. In a similar manner it is found that the complex roots $z_{0, \pm 2}(J)$ and $z_{1, \pm 2}(J)$ lie on a different closed curve $C_{2}$.

It is possible to consider $\zeta_{0}$ to be a parameter for the plane curves $C_{1}$ and $C_{2}$ because for real $J$ the function $\zeta_{0}(J)$ has real values in the interval

$$
\begin{equation*}
-\frac{1}{2}\left[(1-\sqrt{ } 5)+\sqrt{ } 2(5-\sqrt{ } 5)^{1 / 2}\right]<\zeta_{0}(J) \leqslant \frac{1}{2}\left[-(1+\sqrt{ } 5)+\sqrt{ } 2(5+\sqrt{ } 5)^{1 / 2}\right] . \tag{48}
\end{equation*}
$$

We see, therefore, from equation (47) that the closed curves $C_{1}$ and $C_{2}$ can each be described in a piecewise manner by four rational parametric representations. A simplification of this result can be achieved by formally allowing the parameter $\zeta_{0}$ to have any real value. In this manner we find that the curve $C_{1}$ has the single parametric representation

$$
\begin{equation*}
z=z_{\mathrm{c}}\left[\frac{\frac{1}{2}(\sqrt{ } 5-1) \zeta_{0}+\varepsilon}{-\frac{1}{2}(\sqrt{ } 5+1) \zeta_{0}+\varepsilon}\right]^{5} \tag{49}
\end{equation*}
$$

where $-\infty<\zeta_{0}<\infty$. For the curve $C_{2}$ we have the single representation

$$
\begin{equation*}
z=z_{\mathrm{c}}\left[\frac{\frac{1}{2}(\sqrt{ } 5-1) \zeta_{0}+\varepsilon^{2}}{-\frac{1}{2}(\sqrt{ } 5+1) \zeta_{0}+\varepsilon^{2}}\right]^{5} \tag{50}
\end{equation*}
$$

where $-\infty<\zeta_{0}<\infty$. It is clear from equations (49) and (50) that the curves $C_{1}$ and $C_{2}$ are both rational curves.

The cardioid drawn by Wood et al (1989) is part of the curve $C_{1}$ and has the representation (49) with the parameter $\zeta_{0}$ in the intervals

$$
\begin{align*}
& -\frac{1}{4}\left[-(3+\sqrt{ } 5)+\sqrt{ } 6(5+\sqrt{ } 5)^{1 / 2}\right] \leqslant \zeta_{0} \leqslant 0  \tag{51}\\
& \frac{1}{2}(\sqrt{ } 5-1) \leqslant \zeta_{0} \leqslant \frac{1}{4}\left[(3-\sqrt{ } 5)+\sqrt{ } 6(5-\sqrt{ } 5)^{1 / 2}\right] . \tag{52}
\end{align*}
$$

The other curve considered by Wood et al (1989) is part of $C_{2}$ and has the representation (50) with the parameter $\zeta_{0}$ in the intervals

$$
\begin{align*}
& -\frac{1}{2}(\sqrt{ } 5+1) \leqslant \zeta_{0} \leqslant-\frac{1}{4}\left[-(3-\sqrt{ } 5)+\sqrt{ } 6(5-\sqrt{ } 5)^{1 / 2}\right]  \tag{53}\\
& -\frac{1}{4}\left[-(3+\sqrt{ } 5)+\sqrt{ } 6(5+\sqrt{ } 5)^{1 / 2}\right] \leqslant \zeta_{0} \leqslant 0 \tag{54}
\end{align*}
$$

It is interesting to note that the points where the curves $C_{1}$ and $C_{2}$ cross the real $x$ axis are associated with parameter values $\zeta_{0}$ which are zeros of the icosahedral polynomials $\zeta^{5} \Omega_{1}\left(\zeta^{5}\right), \Omega_{2}\left(\zeta^{5}\right)$ and $\Omega_{3}\left(\zeta^{5}\right)$. One can also show that the positions of these crossing points are expressible in terms of the real zeros of the polynomials $\Omega_{j}\left(\zeta^{5}\right)(j=1,2,3)$. For example, the cardioid drawn by Wood et al (1989) crosses the negative $x$ axis at

$$
\begin{equation*}
x=-\zeta_{2}^{5} \tag{55}
\end{equation*}
$$

where

$$
\begin{equation*}
\zeta_{2}=\frac{1}{4}\left[(3-\sqrt{ } 5)+\sqrt{ } 6(5-\sqrt{ } 5)^{1 / 2}\right] \tag{56}
\end{equation*}
$$

is one of the zeros of the polynomial $\Omega_{2}\left(\zeta^{5}\right)$.
It can be shown by equating the imaginary part of equation (14) to zero that the cartesian coordinates $x=x\left(\zeta_{0}\right)$ and $y=y\left(\zeta_{0}\right)$ for both curves $C_{1}$ and $C_{2}$ satisfy a single polynomial equation of the type

$$
\begin{equation*}
P(x, y) \equiv \sum_{n=0}^{10} \sum_{m=0}^{20-2 n} c_{n m} x^{m} y^{2 n}=0 \tag{57}
\end{equation*}
$$

where $c_{n m}$ are non-zero integers. From the parametric equations (49) and (50) one would expect that this equation is reducible to two independent polynomial equations $P_{1}(x, y)=0$ and $P_{2}(x, y)=0$ which are satisfied separately by the curves $C_{1}$ and $C_{2}$. These reduced equations will have coefficients which involve irrational numbers such as $\sqrt{ } 5$.

I am extremely grateful to Dr D W Wood for helpful correspondence and several stimulating discussions on algebraic geometry. I also thank Dr I J Zucker for his expert and generous assistance in the derivation of certain identities involving products of gamma functions. These identities were used to determine the constants $A_{1}, A_{2}, \ldots, A_{9}$ in equations (26), (32) and (38). Finally, I am grateful to Dr J L Martin for his interest in the work, and to Paul Lee for help with computer graphics.

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